

# InP Double Heterojunction Bipolar Transistor for detection above 1 THz

D. Coquillat<sup>1</sup>, V. Nodjiadjim<sup>2</sup>, A. Konczykowska<sup>2</sup>, N. Dyakonova<sup>1</sup>, C. Consejo<sup>1</sup>, S. Ruffenach<sup>1</sup>,  
F. Teppe<sup>1</sup>, M. Riet<sup>2</sup>, A. Muraviev<sup>3</sup>, A. Gutin<sup>3</sup>, M. Shur<sup>3</sup>, J. Godin<sup>2</sup>, W. Knap<sup>1</sup>

<sup>1</sup>Laboratoire Charles Coulomb (L2C), UMR 5221 CNRS-Univ. Montpellier, Montpellier, France

<sup>2</sup>III-V Lab (Bell Labs, TRT and CEA/LETI joint Lab), Route de Nozay, 91460 Marcoussis, France

<sup>3</sup>Rensselaer Polytechnic Institute, Troy, New York 12180, USA

**Abstract**— We evaluate the optical performance of the InP heterojunction bipolar transistors (DHBTs) designed for 100 Gbit/s circuit applications as a room temperature detector operating above 1 THz. They can operate far above the frequencies at which they have gain and can still rectify THz current and voltage.

## I. INTRODUCTION

TERAHERTZ (THz) electromagnetic waves (0.1-10 THz), with the frequencies between microwaves and infrared, have found many applications including THz imaging [1], which has a potential to enable contact-free non-destructive testing of numerous materials, devices, and systems. The THz images taken above 1 THz provide a higher spatial resolution than the images at sub-THz frequencies.

Several groups have also considered the prospect of using a THz carrier frequency for wireless communication to increase the transmitted data rate well above 40 Gbit/s for a single channel [2]. High-data-rate communication was demonstrated very recently using a commercial plasma-wave detector [3]. Error-free communications at data-rates up to 8.2 Gbps were reported at 310 GHz, thus demonstrating that such detectors are competitive for communication applications. However, communications at higher data rates were restricted by the relatively low sensitivities of these detectors. One of the main limitations was their cut-off frequency below 20 GHz.

The THz detectors based on the InP heterojunction bipolar transistors (DHBTs) with high cut-off frequencies are very promising for THz imaging and high-data-rate wireless communication applications. They are favorably compared with the existing detectors based on the Schottky diodes, field effect transistors, or Si heterojunction bipolar transistors [4, 5]. InP DHBTs already serve as a key component of the trans-impedance amplifiers for 100 Gb/s transmission [6].

In this work, we show new results related to the THz detection at room temperature with InP DHBTs in a frequency range above 1 THz.

## II. TECHNOLOGY AND EXPERIMENTAL SET UP

The  $0.7 \times 5 \mu\text{m}^2$  emitter size InP DHBTs were designed for 100 Gbit/s applications. They demonstrated  $f_T$  and  $f_{\text{max}}$  above 320 and 280 GHz respectively. The static current gain was around 40 and the common-emitter breakdown voltage was above 5V. They operated without any spatial coupling antennas. The radiation was coupled to the device only via the contact pads and/or the coupling wires to the device. The photoresponse  $\Delta u$  was measured with a continuous-wave CO<sub>2</sub> -pumped molecular gas laser emitting at 1.40 THz, 1.63 THz, 2.52 THz, and 3.11 THz with maximum output power of  $\sim 100$  mW, 150

mW, 150 mW, and 90 mW, respectively. The radiation was focused on the transistor surface using parabolic mirrors and the radiation intensity was modulated by a chopper at 133 Hz. The focused spot size was close to the diffraction limit. A wire grid polarizer controlled the polarization of the incident THz radiation. The azimuthal angle  $\phi$  (Fig. 1) is the angle between the electric field of the electromagnetic wave and the orientation of the emitter-to-collector direction. The emitter terminal of the device was grounded and the base-emitter bias  $V_{BE}$  was controlled by a Keithley Source Meter. The base-emitter bias and radiation frequency dependence was measured using a lock-in amplifier.

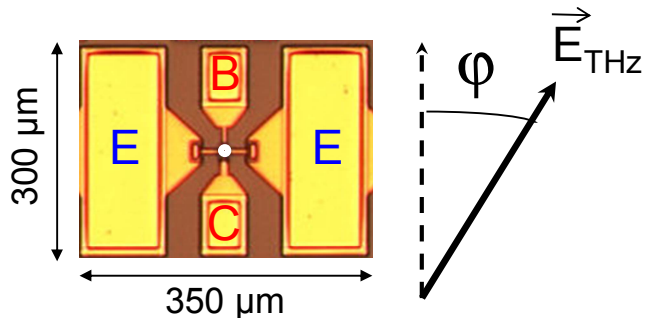


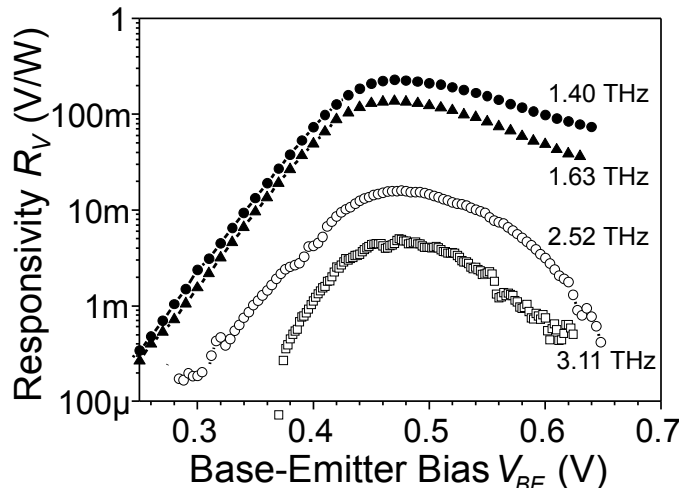
Fig. 1. Image of tested InP DHBT device.  $\phi$  is the polarization angle of incident terahertz radiation.

## III. RESULTS AND DISCUSSION

In Fig. 2 we have plotted the area-normalized responsivity  $R_V$  induced between the collector-emitter electrodes of the  $0.7\text{-}\mu\text{m}$  DHBT as a function of base-emitter bias  $V_{BE}$ , for the four different frequencies. The responsivity  $R_V$  is normalized with regard to the area of the device taken equal to diffraction limited area [7]. In this case,  $R_V$  is defined as the ratio between the voltage  $\Delta u$  induced by the radiation on the detector and the incident power,  $P_t$ .  $R_V = (\Delta u S_t) / (P_t S_a)$ , where  $P_t$  is the total power of the source,  $S_t = \pi d^2 / 4$  is the radiation beam spot area, and  $S_a$  is the active area. Since the area of the DHBT with the contact pads was smaller than the diffraction limited area  $S_\lambda = \lambda^2 / 4$  even for the larger frequency investigated (3.11 THz), the active area was taken equal to  $S_\lambda$ . The values of  $R_V$  used direct reference to laser emitted power without adjustments associated with the factor of coupling efficiency.

The device responsivity rolls off non-linearly with increasing frequency. This observed roll-off can be explained by the combined effects of the capacitance,  $C$ , shunting the active nonlinear device resistance,  $R$ , and parasitic series inductance,  $L$ . At high frequencies the voltage drop  $V_R$  across  $R$  becomes independent of  $R$  and proportional to  $1/(\omega^2 LC)$ . The detector

response is proportional to  $V_R^2$ , and, therefore inversely proportional to  $1/(\omega^4 L^2 C^2)$ .



**Fig. 2.** Area-normalized responsivity  $R_V$  as a function of the base-emitter bias  $V_{BE}$  for frequency of the incident THz radiation of 1.40 THz (●●●), 1.63 THz (▲▲▲), 2.52 THz (○○○), and 3.10 THz (□□□). The polarization of the electric field of the incident THz radiation is perpendicular to the base-to-collector direction.

Moreover, for the polarization of the electric field of the radiation perpendicular to the emitter-to-collector direction, the optimal THz detection bias conditions were achieved for  $V_{BE} \sim 0.48$  V. This value is smaller than the bias voltage that gives the maximum gain in the current voltage characteristic. Above 0.48 V, the responsivity was a decreasing function of  $V_{BE}$ .

#### IV. CONCLUSION

We have demonstrated broadband terahertz photoresponse of InP DHBT and estimated responsivity values. The increase of the responsivity with the *dc* base-emitter bias shows that the detection mechanism is the rectification of the THz radiation coupled to the device due to the nonlinearity of the device characteristics. The roll-off observed with increasing frequency can be explained by the combined effects of the capacitance shunting the active nonlinear device resistance, and parasitic series inductance.

#### ACKNOWLEDGMENTS

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