

# Fast Pyroelectric Response of Semiconducting YBaCuO Detectors with High IR Sensitivity; Development of THz Imaging Arrays

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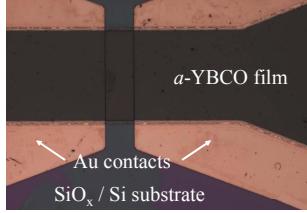
**Abstract**— $\text{YBa}_2\text{Cu}_3\text{O}_{6+x}$  oxides are semiconducting for  $x < 0.5$  and are amorphous ( $a$ -YBCO) at low deposition temperature (100 °C, typically).  $a$ -YBCO also exhibits pyroelectric properties. Metal/ $a$ -YBCO planar micro-structures were processed to fabricate IR sensors. When unbiased, the device exhibited a pyroelectric near-infrared response as  $f^{+1}$  up to about 35 kHz, followed by a slow thermally diffusive decrease as  $f^{-1/2}$  up to about 3 MHz. NEP was measured below 15 pW/Hz $^{1/2}$ , with detectivity  $D^*$  above  $10^9 \text{ cm}\cdot\text{Hz}^{1/2}\cdot\text{W}^{-1}$  in the 500 Hz to 100 kHz range. 32×32 arrays of THz log-periodic planar antennas to be coupled with  $a$ -YBCO patches were also designed.

## I. INTRODUCTION

THE semiconducting phase of  $\text{YBa}_2\text{Cu}_3\text{O}_{6+x}$  (YBCO) with a low oxygen content ( $x < 0.5$ ) offers key advantages for developing uncooled thermal sensors: bolometers [1] or pyroelectric detectors [2]. Moreover, YBCO can be deposited in an amorphous form ( $a$ -YBCO) under low-temperature processing conditions (150 °C, typically), which are compatible with CMOS readout circuitry integration. Recently, very promising performances in terms of detectivity and fast pyroelectric response at infrared wavelengths have been reported with  $a$ -YBCO film detectors [3].

## II. IR RESPONSE RESULTS

In order to confirm the  $a$ -YBCO potential for room temperature imaging, planar structures - 450 nm thick  $a$ -YBCO film connected to two metal pads (see Fig. 1) - were patterned on  $\text{SiO}_x$  (500 nm) /  $p$ -Si (380  $\mu\text{m}$ ) substrates. The device was illuminated with an 850 nm VCSEL diode source modulated in the frequency range  $f = 1 \text{ Hz}$  to 4 MHz. The detector response was readout with a low noise current preamplifier and synchronously detected at  $f$ . The lock-in amplifier was also used to measure the device noise current spectral density in a 1 Hz bandwidth. Besides, noise measurements above 50 kHz were checked with a spectrum analyzer.



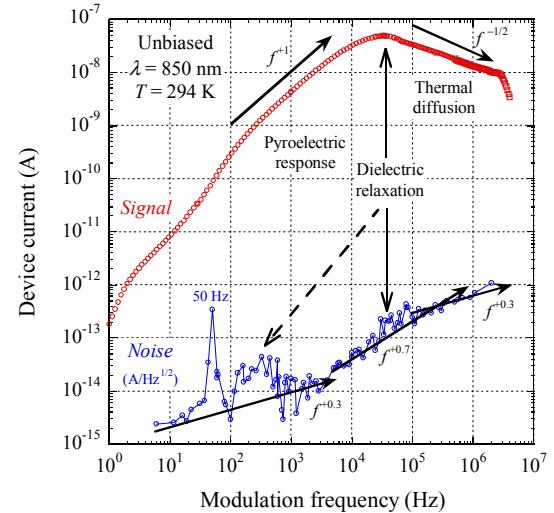
**Fig. 1.** Top view of the as-processed planar device. The  $a$ -YBCO film is 450 nm thick and the sensing micro-bridge dimensions are  $30 \mu\text{m} \times 100 \mu\text{m}$ .

The (unbiased) device current response is shown in Fig. 2. Unbiasing allows eliminating the resistive / bolometric low pass response, and the associated Johnson noise as well. The resulting high pass behavior is typical of a pyroelectric response, i.e., sensitive to the capacitance current (thermally modulated electrical polarization). Such a capacitance current

could be collected because of the  $p$ -doped Si substrate acting as a floating counter-electrode [3]. For this high-pass  $f^{+1}$  pyroelectric behavior, the frequency cut-off at ~35 kHz can be related with dielectric relaxation in the  $a$ -YBCO material [3], [4]. The higher frequency low-pass response – close to  $f^{-1/2}$  behavior typical of thermal diffusion across the substrate – is limited by the current preamplifier at ~3 MHz.

The device current noise ( $i_N$ ) spectrum, as shown in Fig. 2, exhibits a rather unusual “blue noise” behavior, which can be thought as reflecting  $a$ -YBCO AC conductivity ( $\sigma_{AC}$ ) dependence with frequency. As a consequence of the fluctuation-dissipation theorem, one can expect that the noise power density  $\langle i_N^2 \rangle(f) \propto \sigma_{AC}(f)$ . The main features observed are:  $i_N \propto f^{0.3}$  below a few kHz (rough trend), then a significant rise  $i_N \propto f^{0.7}$  between a few kHz and 100 kHz (the dielectric relaxation region), and  $i_N \propto f^{0.3}$  above 100 kHz.

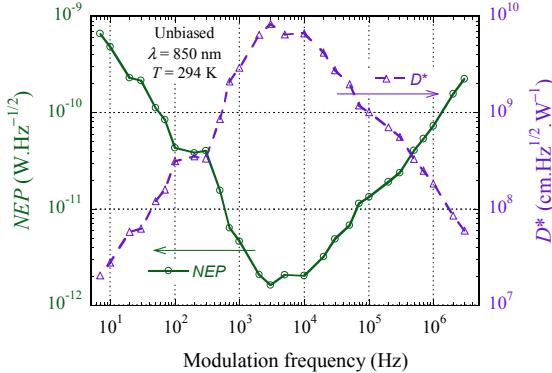
From the correlated barrier hopping (CBH) model, it follows that  $\sigma_{AC}(f) \propto f^n$  with  $n = 1 - 6k_B T/E_b$ , where  $E_b$  is the energy barrier for the hopping process that can be deduced from optical measurements [5]. We thus measured  $E_b = 0.46 \text{ eV}$ , leading to  $n = 0.67$  at room temperature, whereas we directly measured  $\sigma_{AC} \propto f^{0.71}$  for  $10 \text{ kHz} < f < 100 \text{ kHz}$  as explained in [4]. We should therefore expect  $i_N \propto f^{0.35}$  (i.e.,  $\sim f^{0.71/2}$ ) in this latter frequency range from noise measurements. The remaining 0.3 to 0.35 exponent as exhibited in Fig. 2 in the whole frequency range (assuming additivity of noise power), may be attributed to a hopping mechanism across a contact barrier, with hopping energy around 0.4 to 0.5 eV.



**Fig. 2.** Typical room temperature NIR response (top) and noise spectrum (bottom) of the unbiased  $a$ -YBCO planar device shown in Fig. 1.

Two extra features can be noticed in the noise spectrum, in the form of “bumps” arising at a few hundred Hz and at a few tens of kHz. These bumps can be related to fluctuations in the dielectric response, which includes a low frequency relaxation and a high frequency relaxation, as previously reported [6].

Noise measurements allowed to extract the NEP and detectivity  $D^*$  of the device. As shown in Fig. 3, NEP remains below 15 pW/Hz $^{1/2}$ , and  $D^*$  above 10 $^9$  cm·Hz $^{1/2}\cdot\text{W}^{-1}$ , in the 500 Hz to 100 kHz range. With minimum NEP  $\approx$  2 pW/Hz $^{1/2}$ , and maximum  $D^* \approx 7 \times 10^9$  cm·Hz $^{1/2}\cdot\text{W}^{-1}$  (theoretical room temperature limit  $\sim 2 \times 10^{10}$  cm·Hz $^{1/2}\cdot\text{W}^{-1}$ ), this pyroelectric detector is quite competitive compared to other reported NIR thermal detectors (see Table 1). The very short response time  $\tau \approx 2 \mu\text{s}$  should also be noticed.



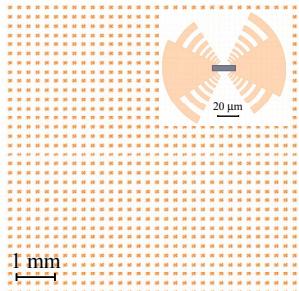
**Fig. 3.** For the unbiased  $a$ -YBCO planar device, NIR noise equivalent power and detectivity as a function of modulation frequency.

**Table 1.** For planar structures based on amorphous and/or oxide materials, room temperature performance for NIR thermal detection.

Material ref.	$f$ (Hz)	NEP (W/Hz $^{1/2}$ )	$D^*$ (cm·Hz $^{1/2}/\text{W}$ )	$\tau$ (s)
$a$ -YBCO this work	200	3.8×10 $^{-11}$	3.5×10 $^8$	
	1×10 $^4$	2×10 $^{-12}$	6.6×10 $^9$	1.9×10 $^{-6}$
	1×10 $^5$	1.3×10 $^{-11}$	1.0×10 $^9$	
$a$ -YBCO [7]	30	4.1×10 $^{-11}$	1.2×10 $^8$	5.7×10 $^{-3}$
$a$ -Si [8]	200	8×10 $^{-12}$	1.3×10 $^8$	2.1×10 $^{-3}$
$a$ -SiGe [9]	n.a.	n.a.	2.5×10 $^9$	125×10 $^{-3}$
$a$ -SiGeO [10]	250	n.a.	6.7×10 $^8$	13×10 $^{-3}$
VWOx [11]	15	5.4×10 $^{-10}$	1.1×10 $^7$	7.2×10 $^{-4}$
MCO [12]	30	3.7×10 $^{-10}$	6×10 $^7$	19×10 $^{-3}$

### III. THZ DEVELOPMENT

Due to  $a$ -YBCO reflectivity in the THz [13], coupling of the radiation to the sensing area is being performed using a planar antenna. We have considered the design of antenna arrays; an example of 32×32 array is shown in Fig. 4. A major outcome is the overall array size, according to the criteria of both THz bandwidth and inter-pixel cross talk. Another aim is to test the antenna impedance matching with the  $a$ -YBCO patch using FTIR spectroscopy measurements on such dense arrays of adequately loaded micro-antennas.



**Fig. 4.** THz pixel array design. Inset: detail of a loaded log-periodic antenna.

### IV. SUMMARY

Simple planar detectors structures based on  $a$ -YBCO films have been designed and tested. The optical response of the device can be interpreted in terms of a pyroelectric capacitive current. Competitive NEP and detectivity are exhibited in the near-infrared, together with fast response in the microsecond range. THz imaging arrays using these detectors are further considered.

### V. ACKNOWLEDGMENTS

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